

NEUTRON GRATING INTERFEROMETRY

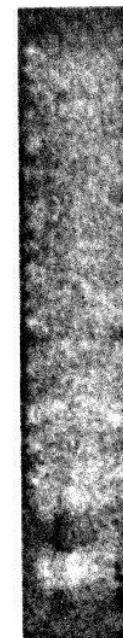
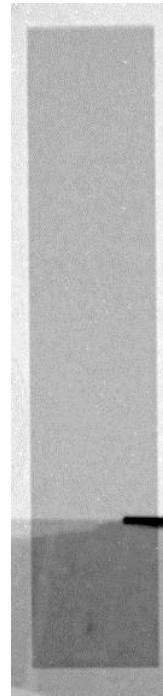
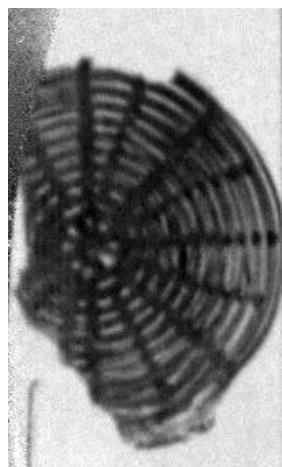
PART II

QUANTITATIVE DFI

Alexander Backs
alexander.backs@frm2.tum.de

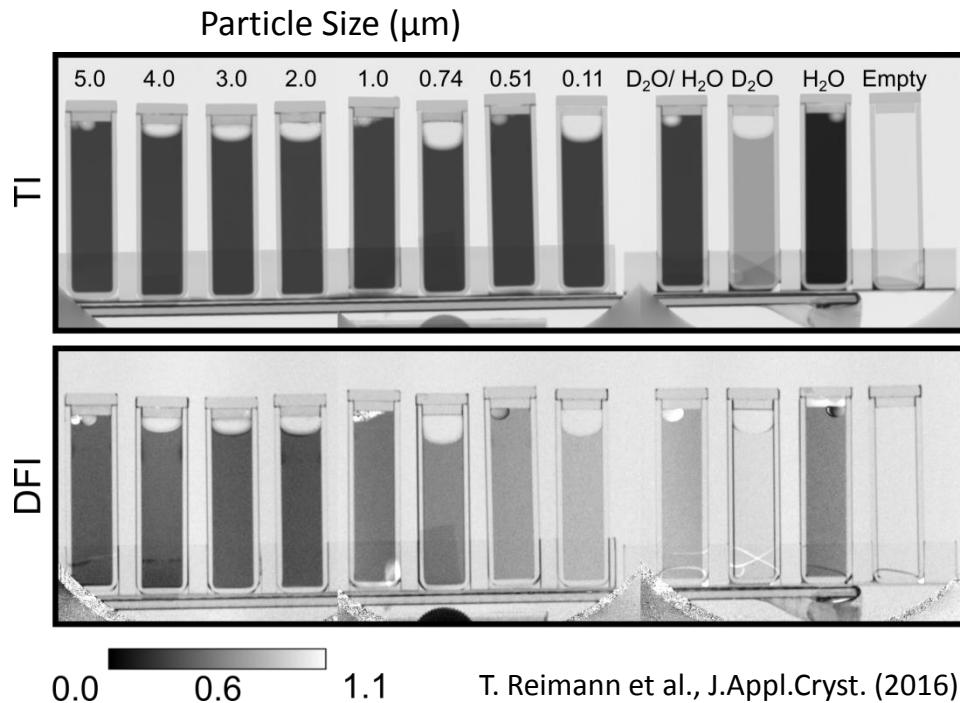
nGI reveals scattering

- from small structures ($\sim \mu\text{m}$)
- under very small angles



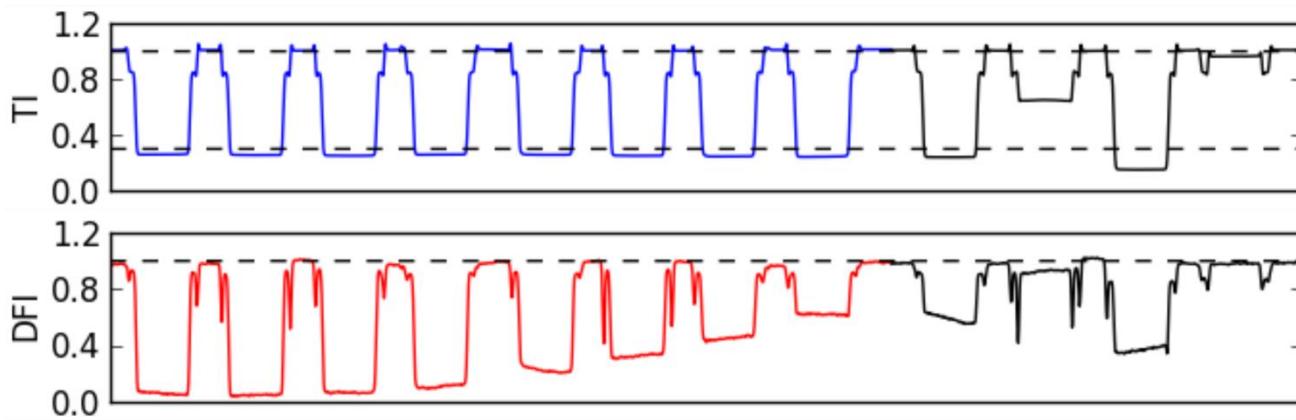
What does that mean exactly ?

Quantitative DFI: Size Matters



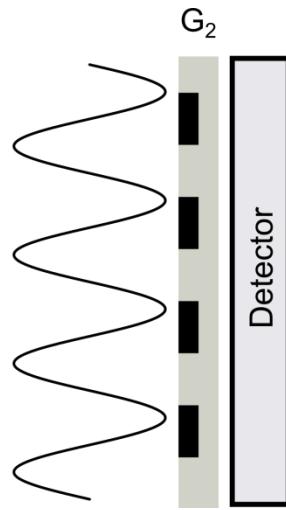
Polyesterene Colloids in Solution

- varying particle diameter
 - constant volume fraction
- ↓
- constant absorbtion
 - varying DFI value



Origin of the DFI

Model Case: One Single Scattering Angle

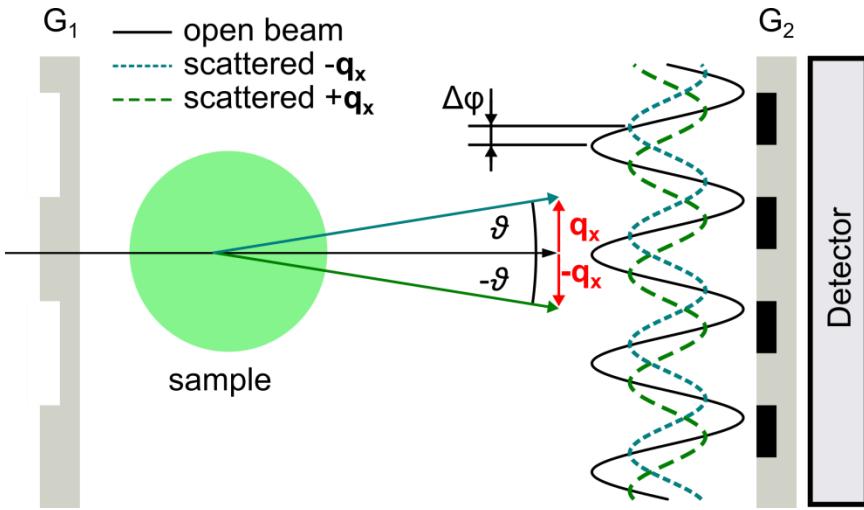


without sample:

$$V_{ob} = \cos(x)$$

Origin of the DFI

Model Case: One Single Scattering Angle



without sample:

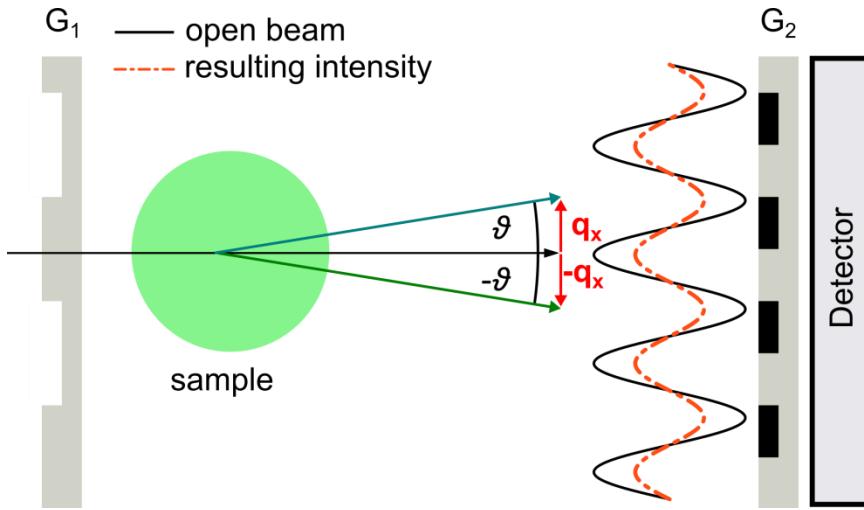
$$V_{ob} = \cos(x)$$

with sample:

$$V_s = \frac{1}{2} \cos(x + \vartheta) + \frac{1}{2} \cos(x - \vartheta)$$

Origin of the DFI

Model Case: One Single Scattering Angle



without sample:

$$V_{ob} = \cos(x)$$

with sample:

$$V_s = \frac{1}{2} \cos(x + \vartheta) + \frac{1}{2} \cos(x - \vartheta)$$

$$V_s = \cos(x) \cos(\xi_{GI} q_x)$$

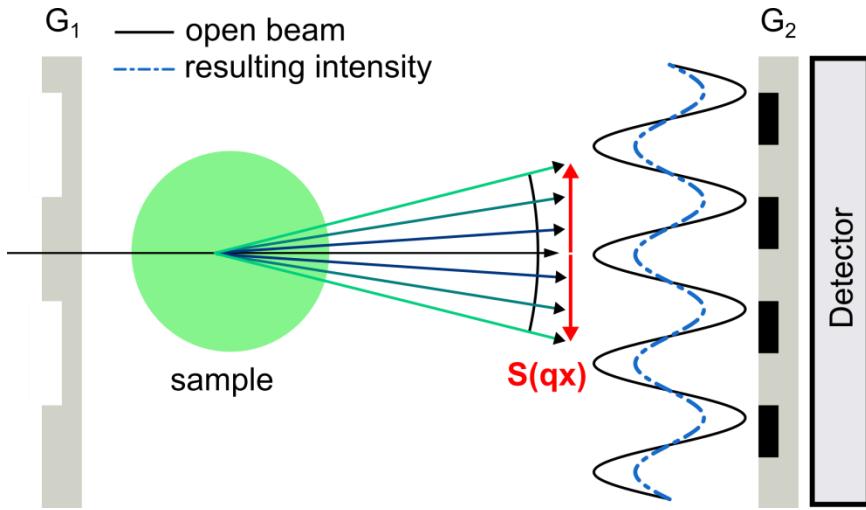
$$\vartheta = \xi_{GI} q_x$$

$$\xi_{GI} = \lambda \frac{L_{s,eff}}{p_2}$$

$$DFI = \frac{V_s}{V_{ob}} = \cos(\xi_{GI} q_x)$$

Origin of the DFI

Real Case: Distribution of Scattering Angles

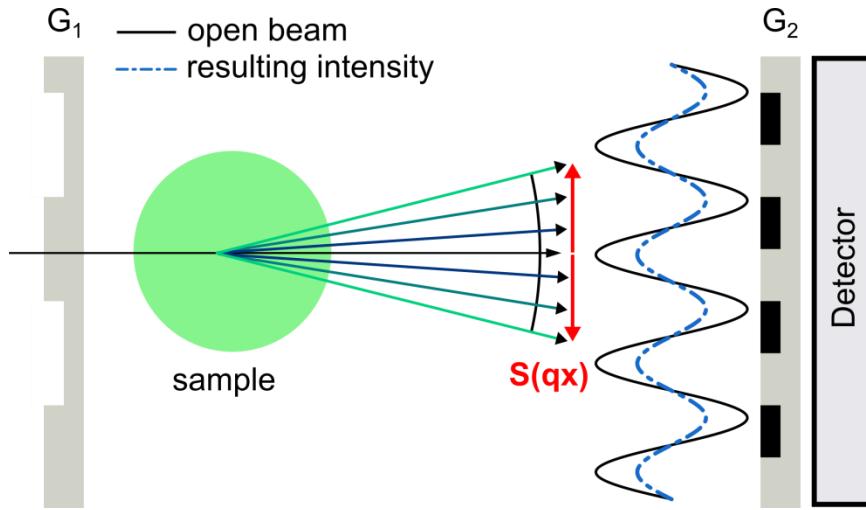


$$DFI = \cos(\xi_{GI} q_x)$$

$$DFI = \int S(q_x) \cos(\xi_{GI} q_x) dq_x$$

Origin of the DFI

Real Case: Distribution of Scattering Angles



$$DFI = \cos(\xi_{GI} q_x)$$

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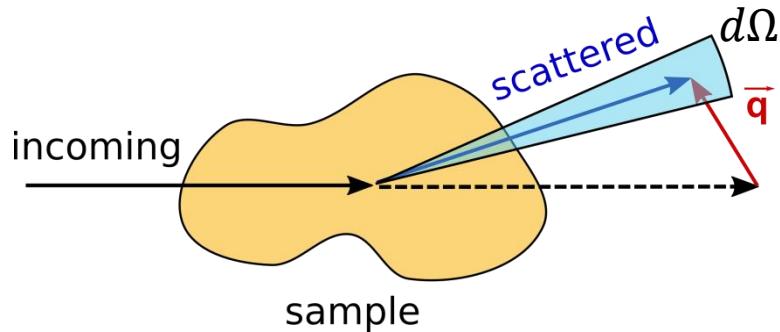
$S(q_x)$

“probability of scattering a neutron with a certain momentum transfer q_x “

What defines $S(q_x)$?

$S(q_x)$: A Bit of Scattering Theory

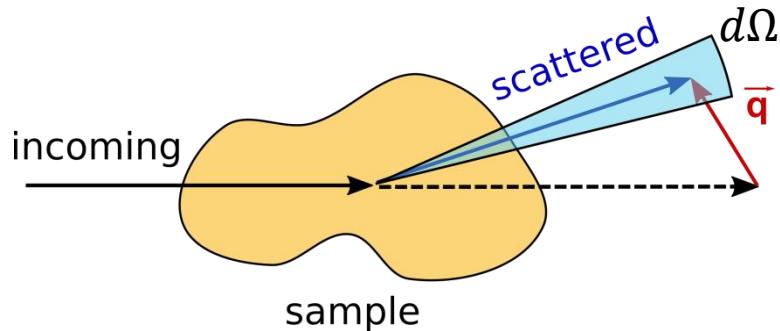
Reciprocal Space: the differential scattering cross section
(looking at momentum)



$$\frac{d\sigma(q)}{d\Omega} = \frac{\text{neutrons scattered into } d\Omega}{\text{incoming neutrons per unit area}}$$

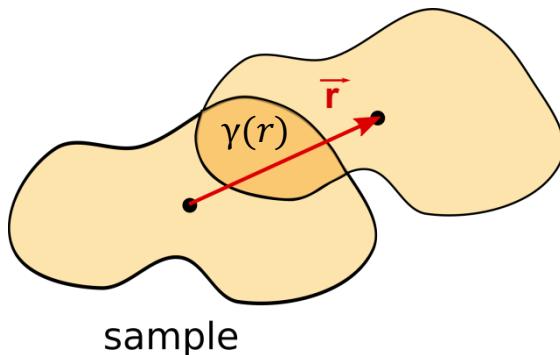
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Real space: the pair correlation function
(looking at space coordinates)

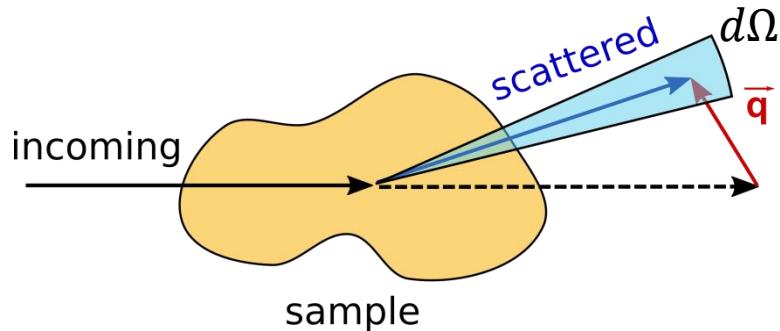


$$\gamma(r) = \int_V \Delta\rho(R)\Delta\rho(R + r)dR$$

Distribution of scattering strength

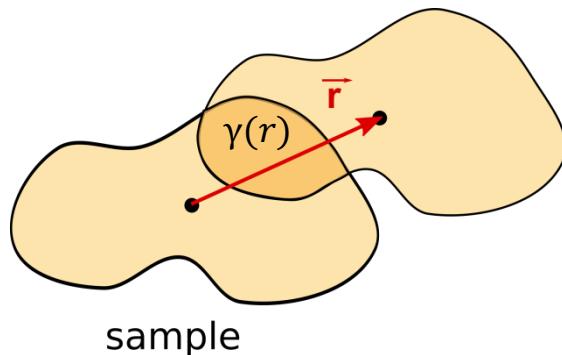
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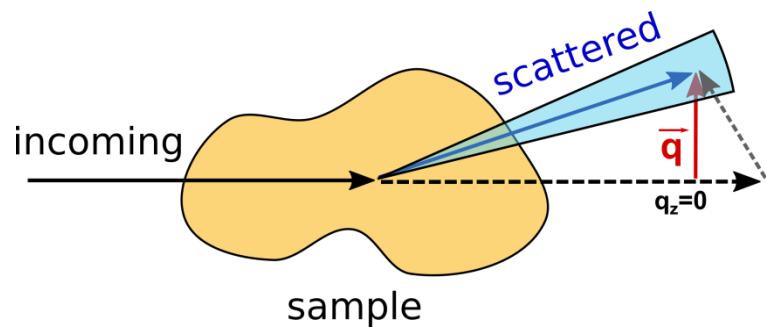
fourier transformation

$$\gamma(r) = \int_V \Delta\rho(R)\Delta\rho(R + r)dR$$

Distribution of scattering strength

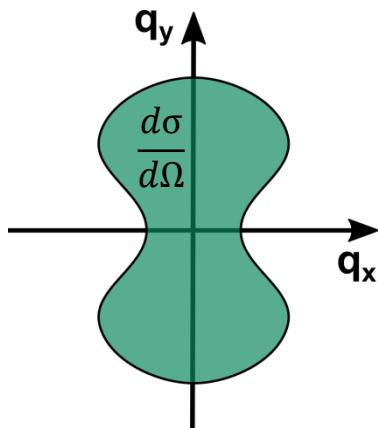
$S(q_x)$: Reducing the Dimensions

Step 1: no scattering along the beam direction $q_z = 0$



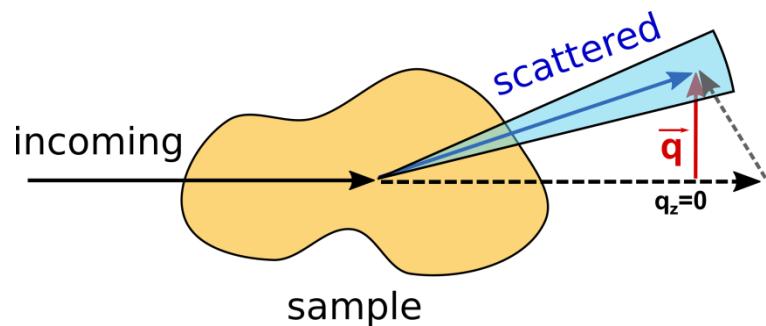
$$\frac{d\sigma(q)}{d\Omega} \rightarrow \frac{d\sigma(q_x, q_y, 0)}{d\Omega}$$

$$\Rightarrow G(x, y) = \int \gamma(x, y, z) dz$$



$S(q_x)$: Reducing the Dimensions

Step 1: no scattering along the beam direction $q_z = 0$

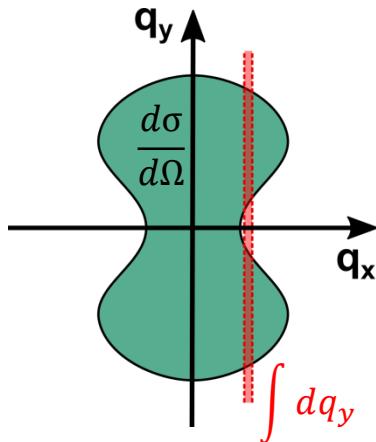


$$\frac{d\sigma(q)}{d\Omega} \rightarrow \frac{d\sigma(q_x, q_y, 0)}{d\Omega}$$

$$\Rightarrow G(x, y) = \int \gamma(x, y, z) dz$$

Step 2: projection along the grating lines

$$q_y = -\infty \rightarrow +\infty$$



$$\left(\frac{d\sigma(q_x)}{d\Omega} \right)_{slit} = \int \frac{d\sigma(q_x, q_y, 0)}{d\Omega} dq_y$$

$$\Rightarrow G(x, y) \rightarrow G(x, 0)$$

Quantitative DFI: The Formula

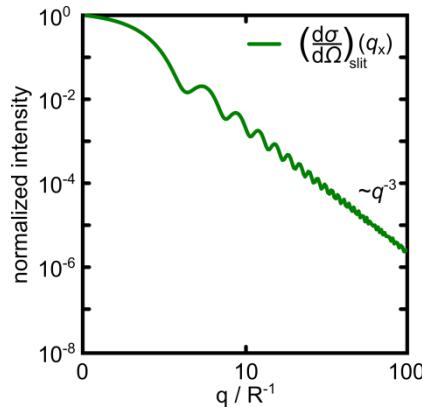
Reciprocal Space:

$$S(q_x) \approx \left(\frac{d\sigma(q_x)}{d\Omega} \right)_{slit}$$

$$\begin{aligned} q_z &= 0 \\ q_y &= -\infty \rightarrow +\infty \end{aligned}$$

↓

$$DFI = \int S(q_x) \cos(\xi_{GI} q_x) dq_x$$



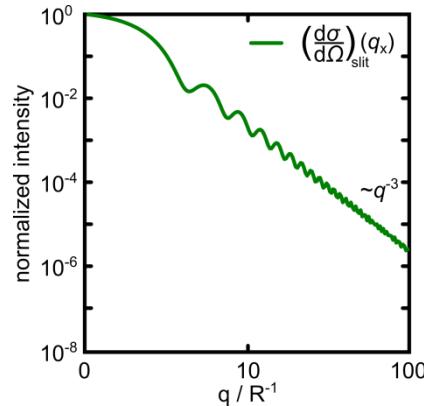
Quantitative DFI: The Formula

Reciprocal Space:

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 $q_y = -\infty \rightarrow +\infty$

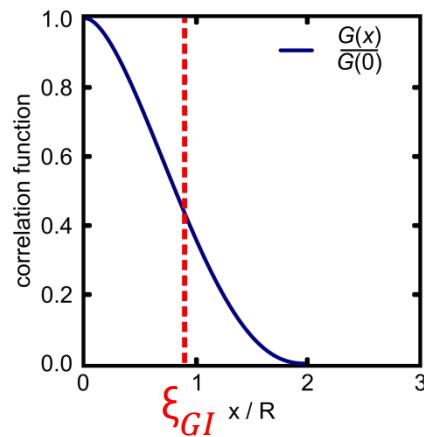
$$DFI = \int S(q_x) \cos(\xi_{GI} q_x) dq_x$$



Real Space:

$$DFI = \exp \left[\Sigma t \left(\frac{G(\xi_{GI})}{G(0)} - 1 \right) \right]$$

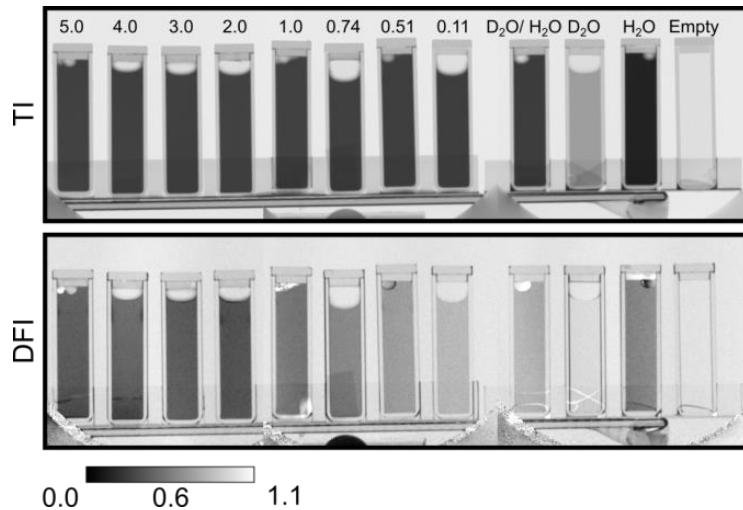
Total scattering crosssection	$\Sigma = (\Delta\rho)^2 \lambda^2 \Phi_v \frac{G(0)}{\gamma(0)}$
sample thickness	t
nGI correlation length	ξ_{GI}



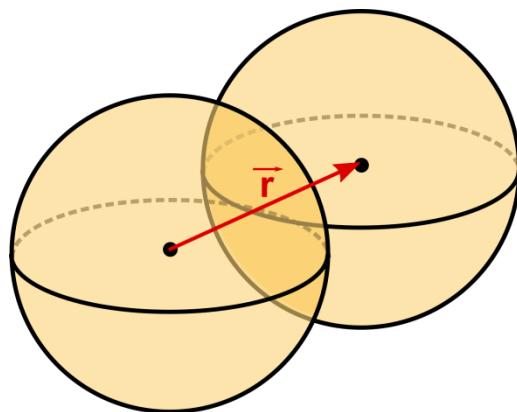
T. Reimann, Ph.D. Thesis (2016)

A Simple Example: Hard Spheres

Polyesterene Colloids in Solution

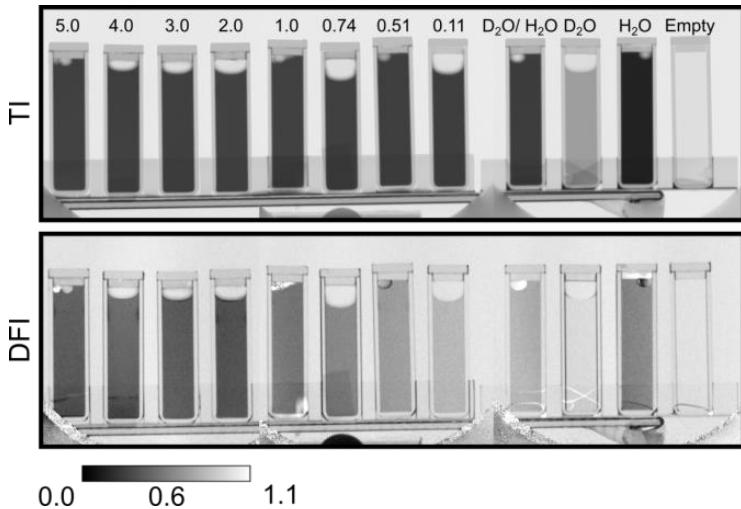


Pair Correlation Function

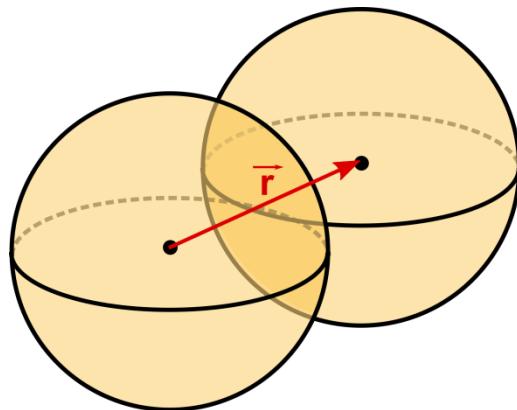


A Simple Example: Hard Spheres

Polyesterene Colloids in Solution



Pair Correlation Function



calculate:

$$\gamma(x, y, z)$$



fourier transform:

$$\frac{d\sigma(q_x, q_y, q_z)}{d\Omega}$$

$$G(x, y)$$



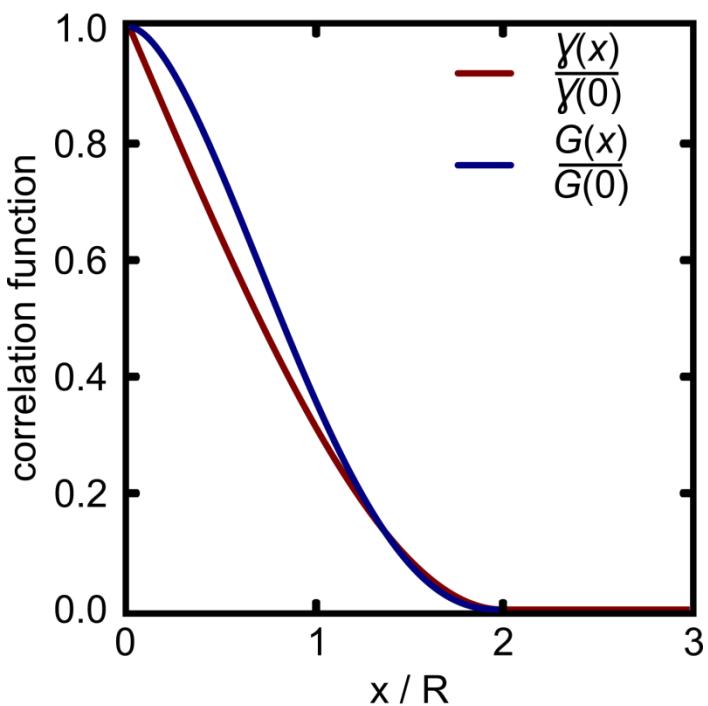
$$G(x, 0)$$



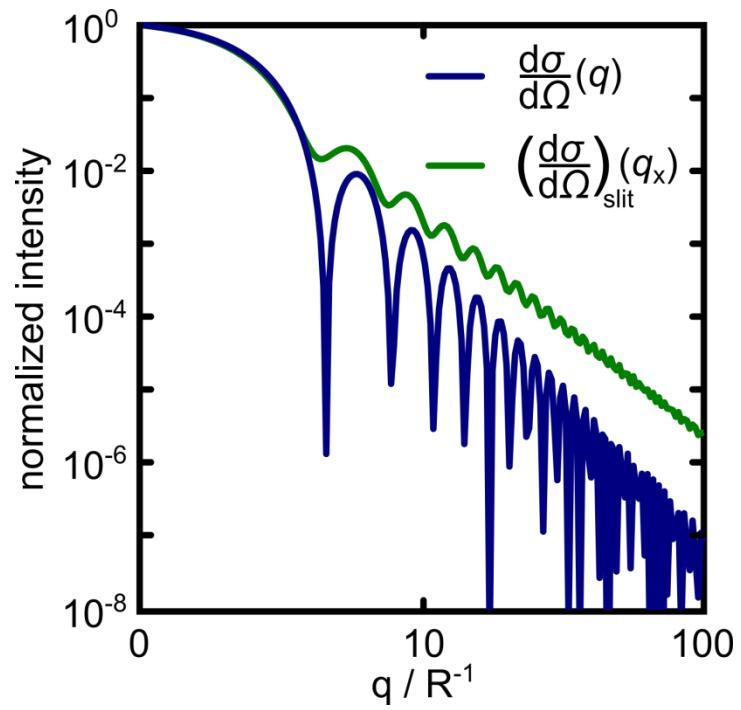
Real Space

Reciprocal Space

Real Space:
Pair Correlation function



Reciprocal Space:
Scattering cross section



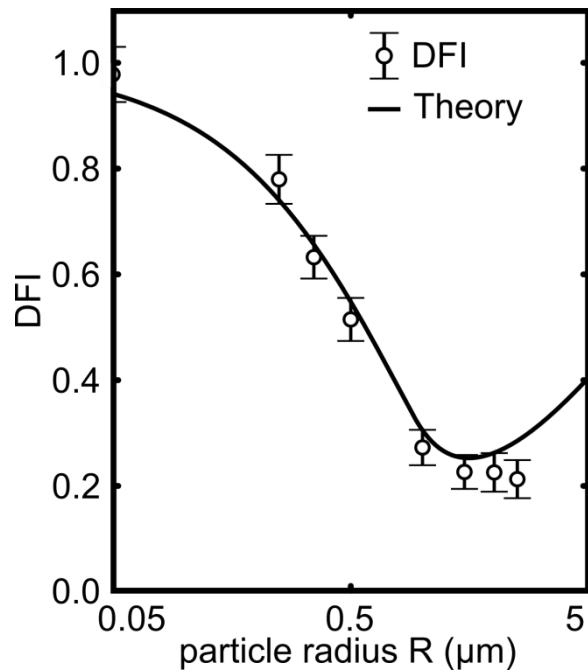
T. Reimann et al., J.Appl.Cryst. (2016)

Polystyrene Colloids in Water

$$DFI = \exp \left[\Sigma t \left(\frac{G(\xi_{GI})}{G(0)} - 1 \right) \right]$$

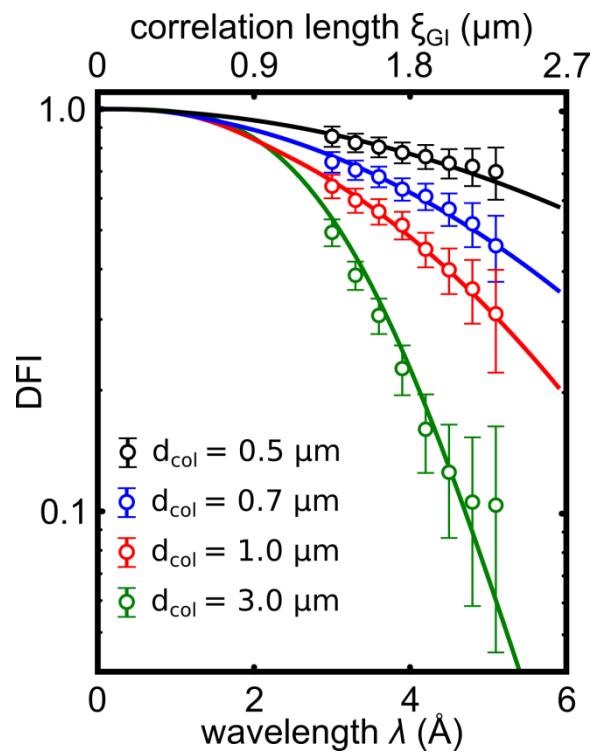
R dependence:

- $G = G(R)$
- $\Sigma = \Sigma(R)$



λ dependence:

- $\xi_{GI} = \xi_{GI}(\lambda)$
- $\Sigma = \Sigma(\lambda)$



T. Reimann et al., J.Appl.Cryst. (2016)

DFI depends only on **one** value of the pair correlation function $G(\xi_{GI})$

More information can be gained by:

Changing G



Different particle sizes



Only possible in rare cases

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Changing ξ_{GI}



Change the setup
or
Sample position



Various geometric
constraints
Tedious calibrations

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Various geometric
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Tedious calibrations

Vary the wavelength

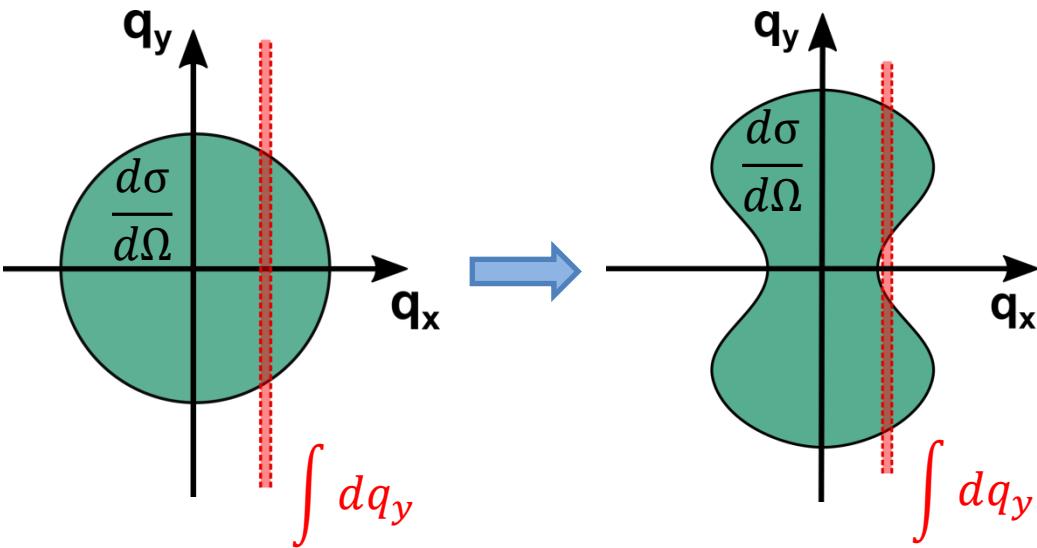


Only a small range
is accessible

Inhomogeneous Scattering

homogenous

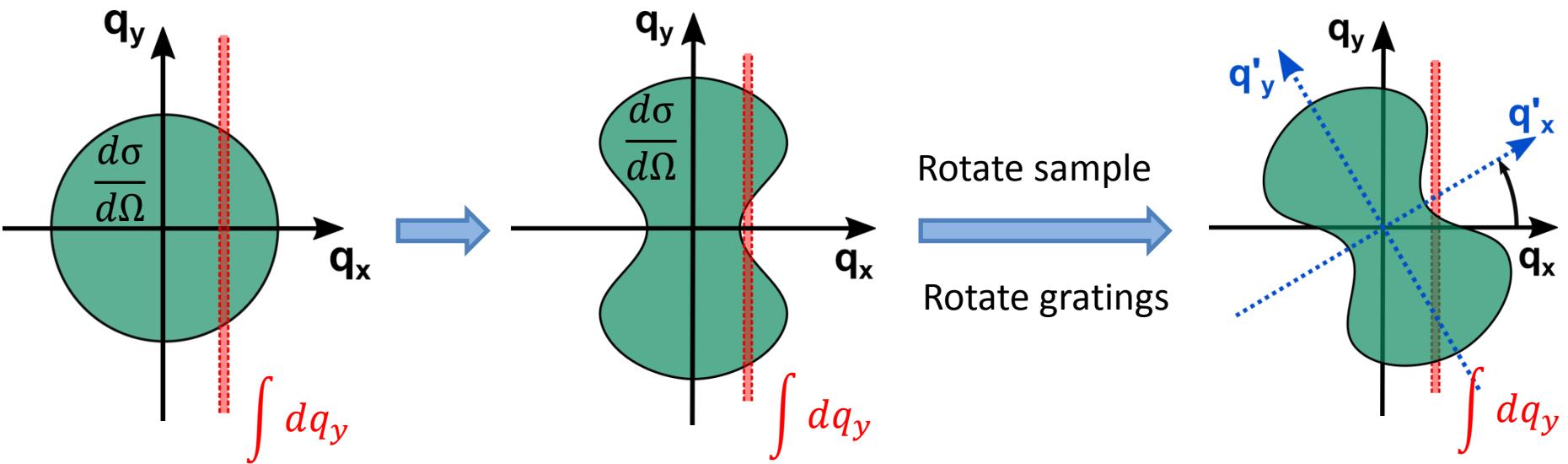
inhomogenous



Inhomogeneous Scattering

homogenous

inhomogenous

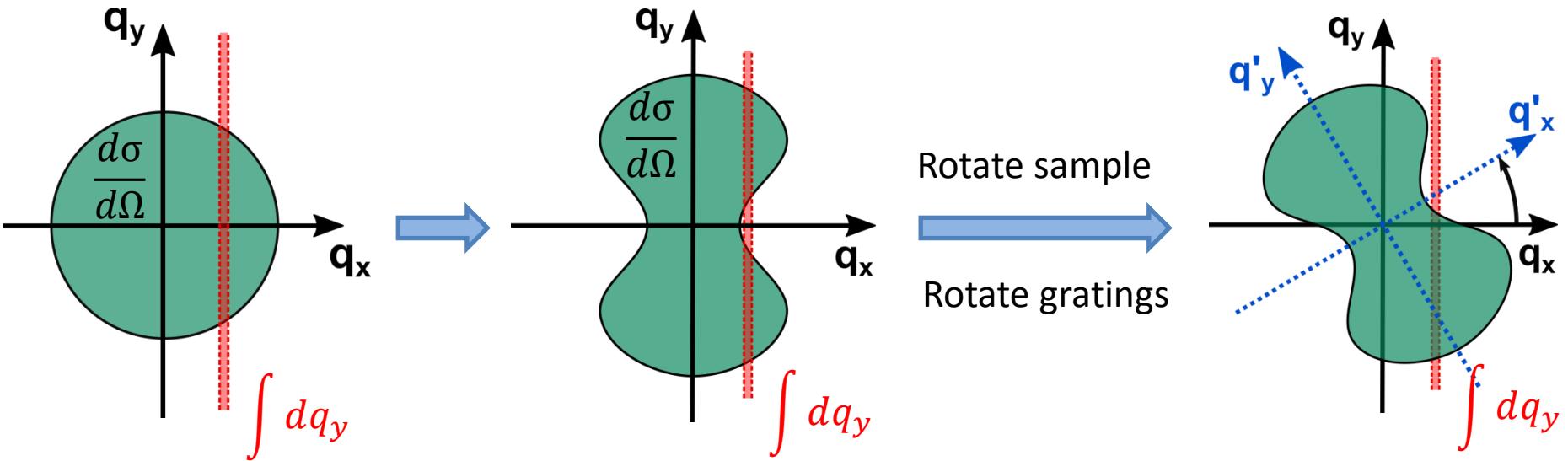


$\left(\frac{d\sigma(q_x)}{d\Omega}\right)_{slit}$ depends on the relative orientation of sample and gratings

Inhomogeneous Scattering

homogenous

inhomogenous



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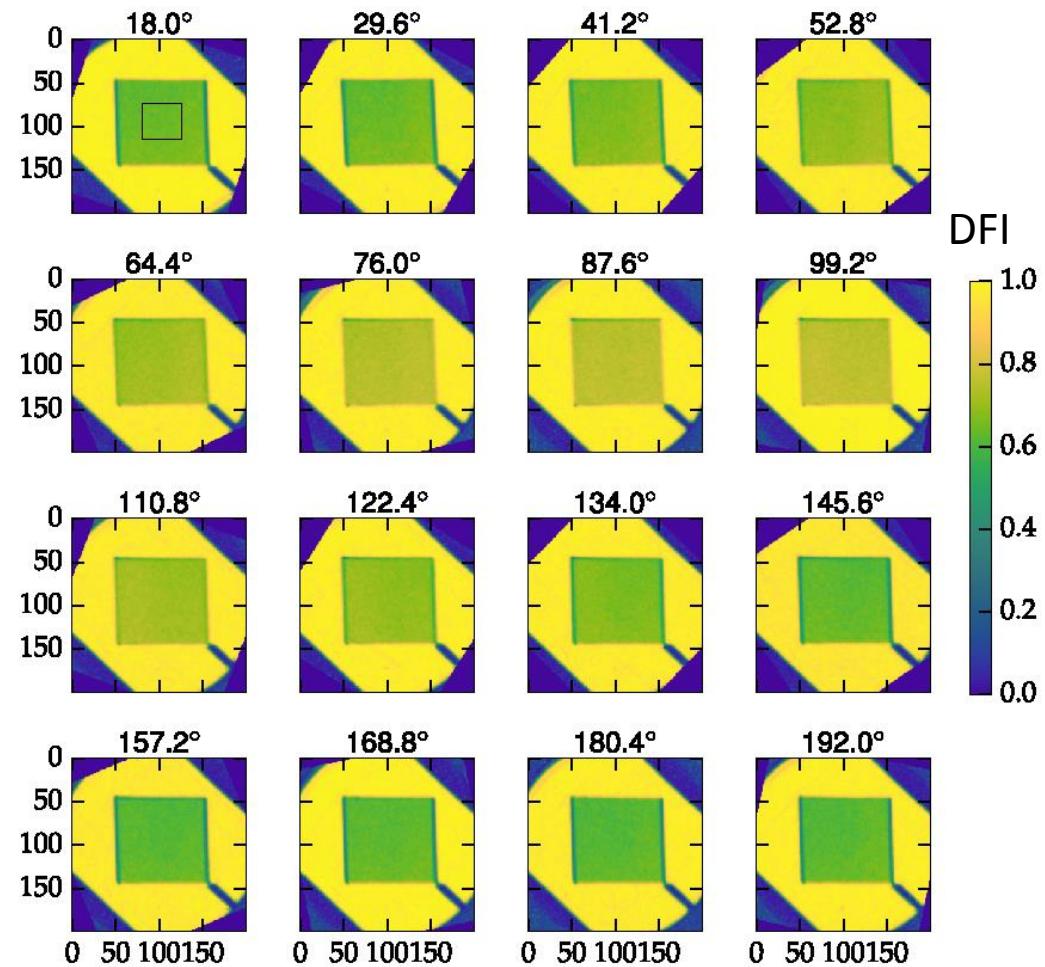
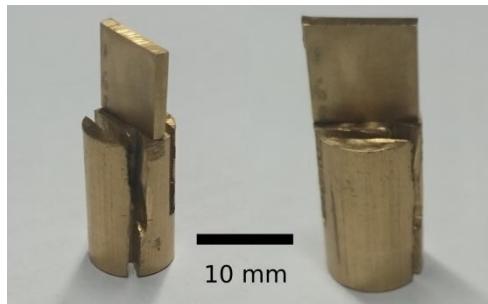
$$\rightarrow DFI(\omega) = \exp \left[\Sigma t \left(\frac{G(\xi_{GI} \cos(\omega), -\xi_{GI} \sin(\omega))}{G(0)} - 1 \right) \right]$$

Anisotropic scattering in Brass

Extrusion Moulded Brass:

anisotropic

- production
- crystallites
- scattering



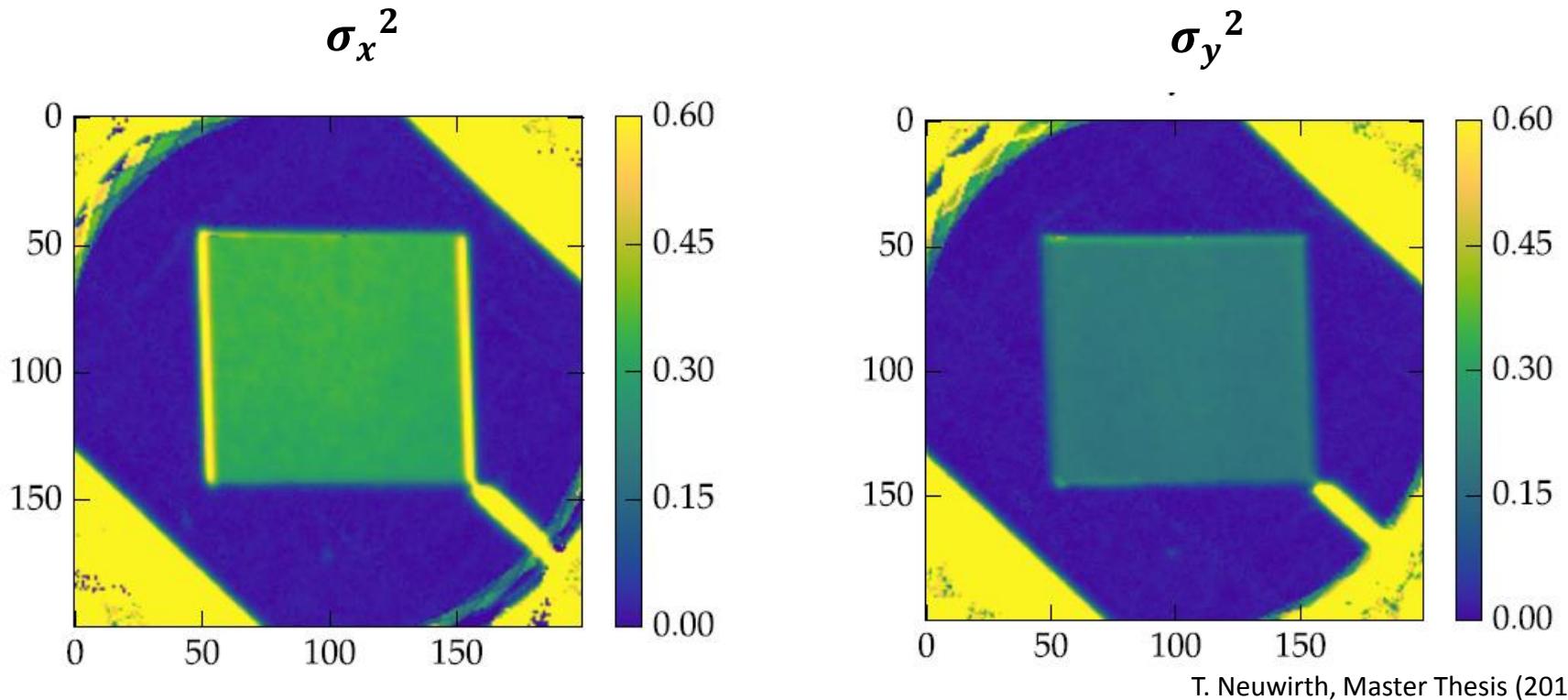
T. Neuwirth, Master Thesis (2017)

Anisotropic scattering in Brass

Different distribution of crystallite size in x and y direction

→ Bi gaussian scattering cross section

$$DFI(\omega) = \exp \left[-\frac{\xi_{GI}}{2} (\sigma_x^2 + (\sigma_y^2 - \sigma_x^2) \sin^2(\omega - \varphi)) \right] \quad \sigma_{an} = \sigma_x^2 - \sigma_y^2$$



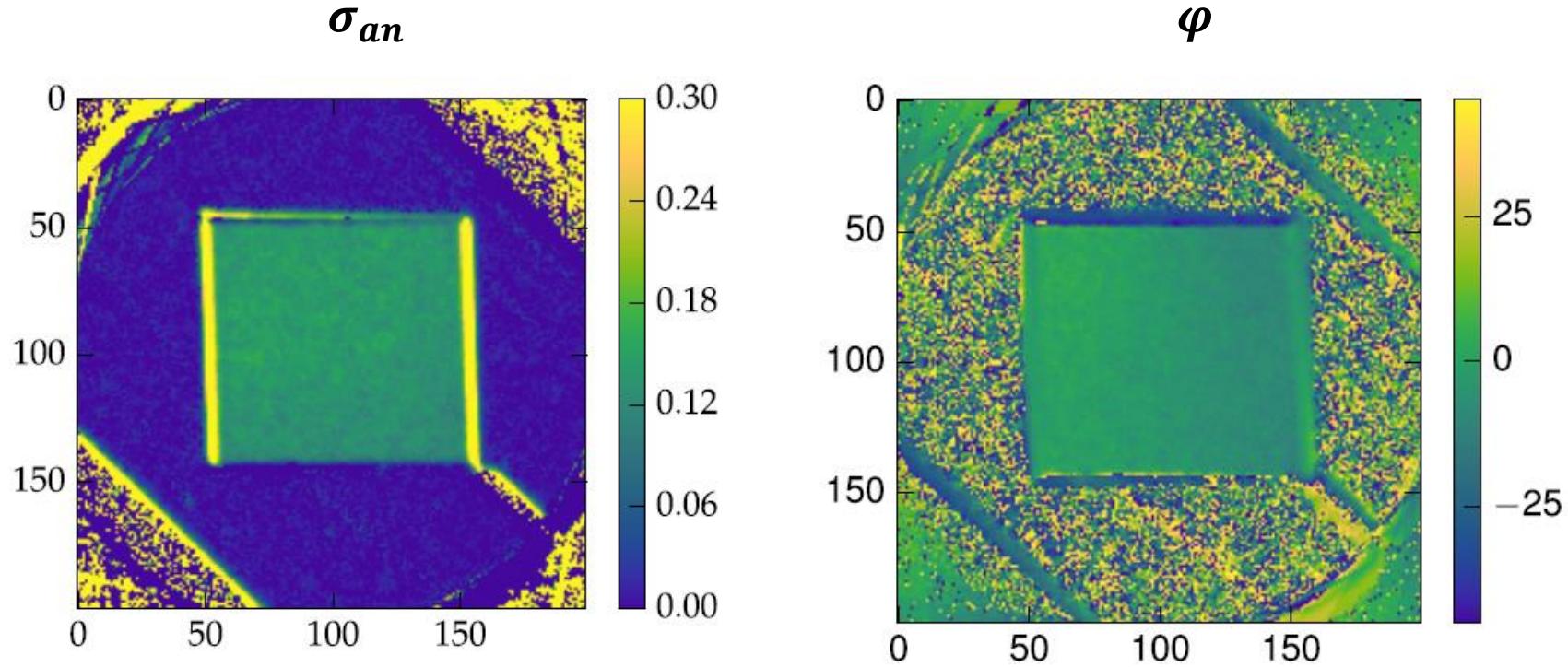
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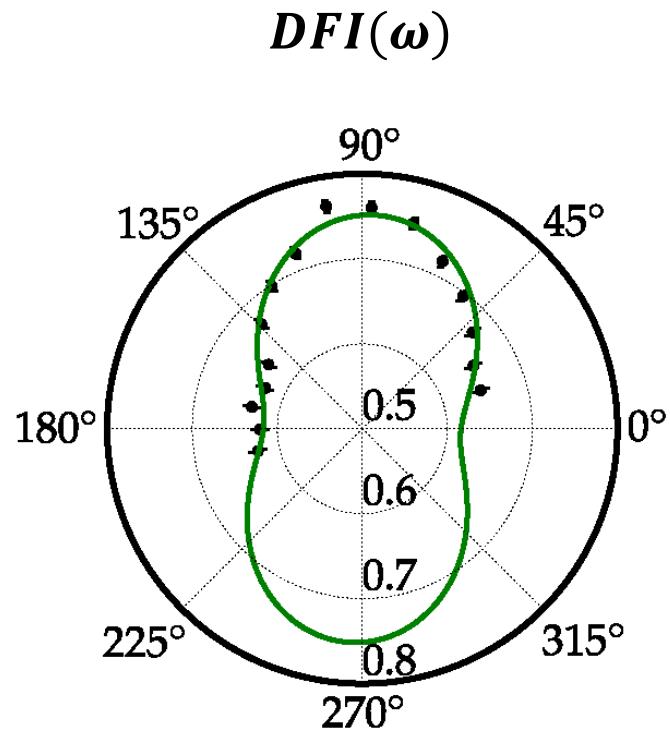
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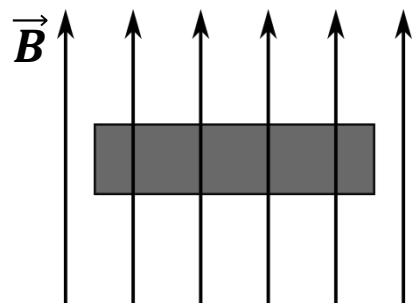
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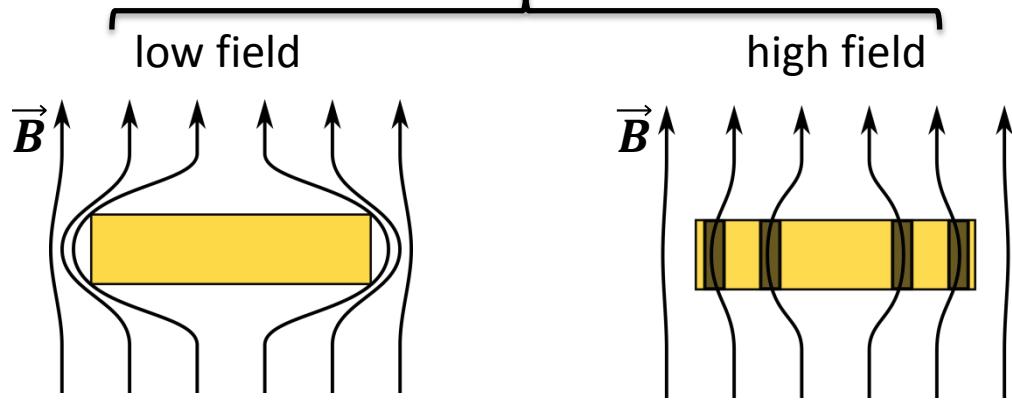
T. Neuwirth, Master Thesis (2017)

Magnetic Domains in Superconductors

Normal Conductor



Superconductor

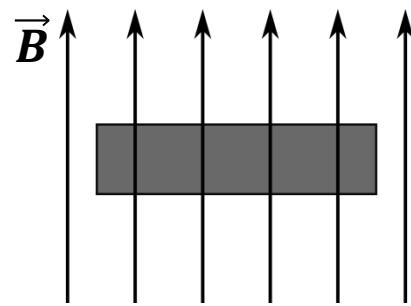


Meissner state

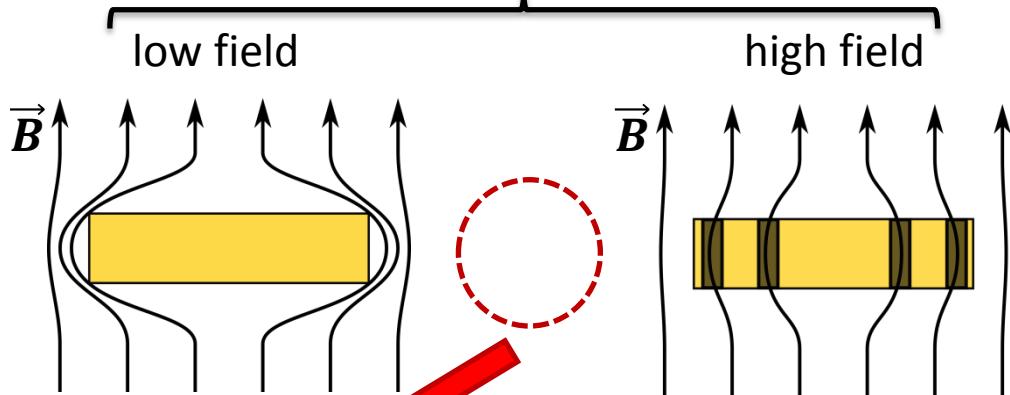
Schubnikov state

Magnetic Domains in Superconductors

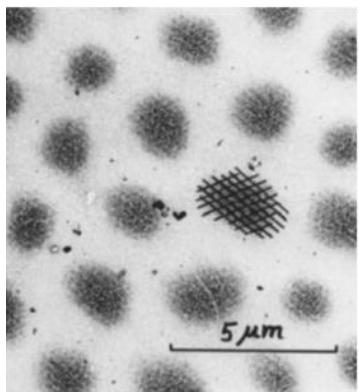
Normal Conductor



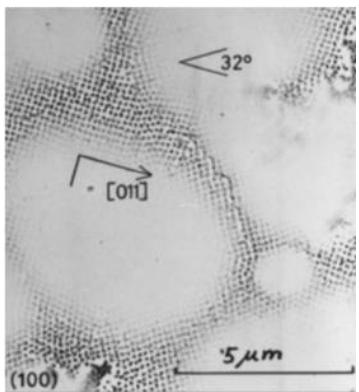
Superconductor



intermediate mixed state bubble structure

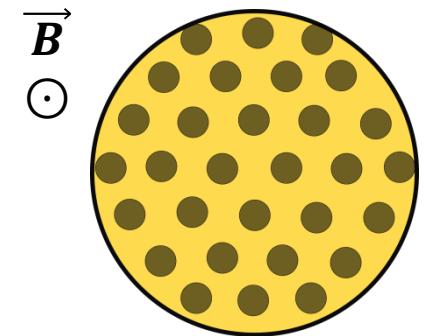
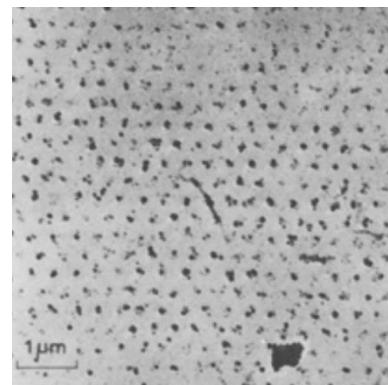


laminar structure



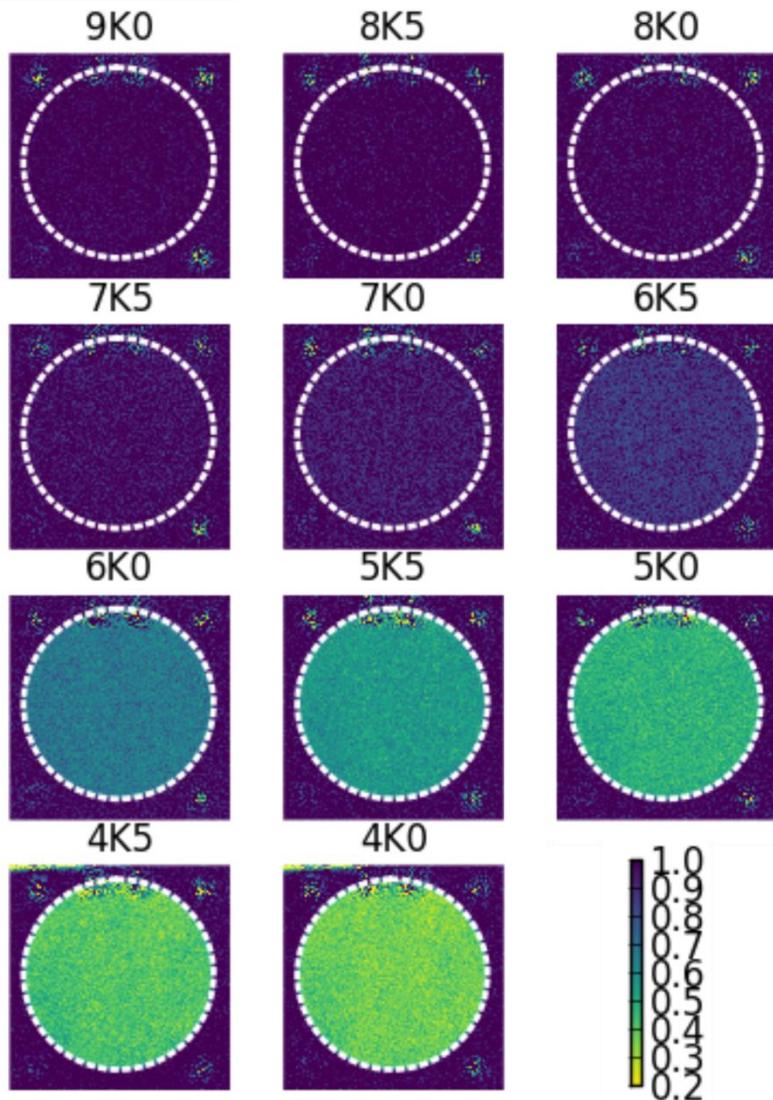
E. H. Brandt & M. P. Das, J. Supercond. Nov. Magn. (2011)

flux line lattice

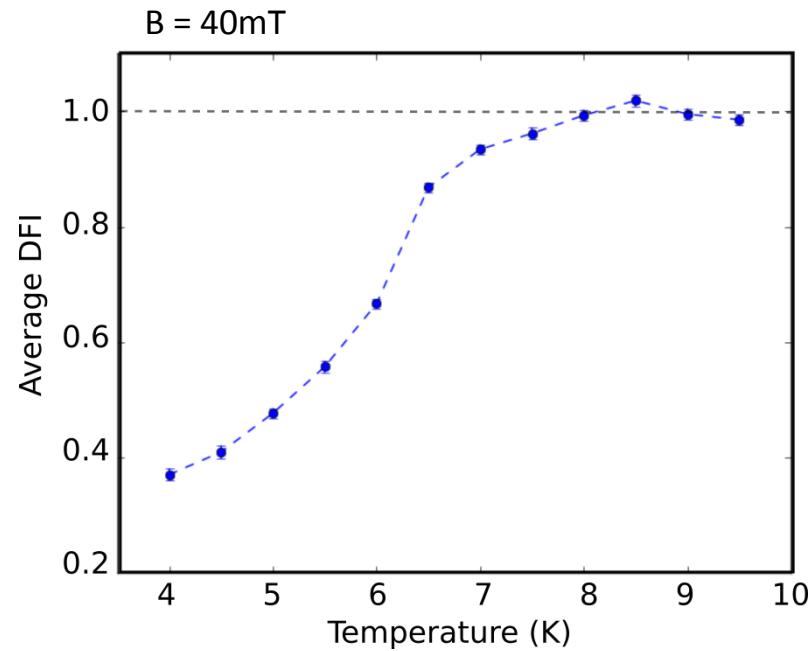


U. Essmann & H. Träuble, Phys. Lett. A (1967)

The Intermediate Mixed State in Niobium



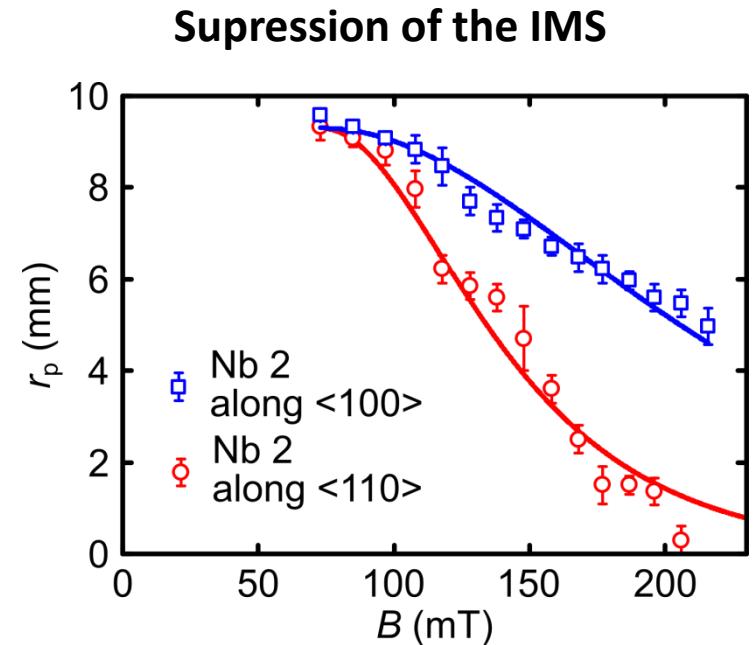
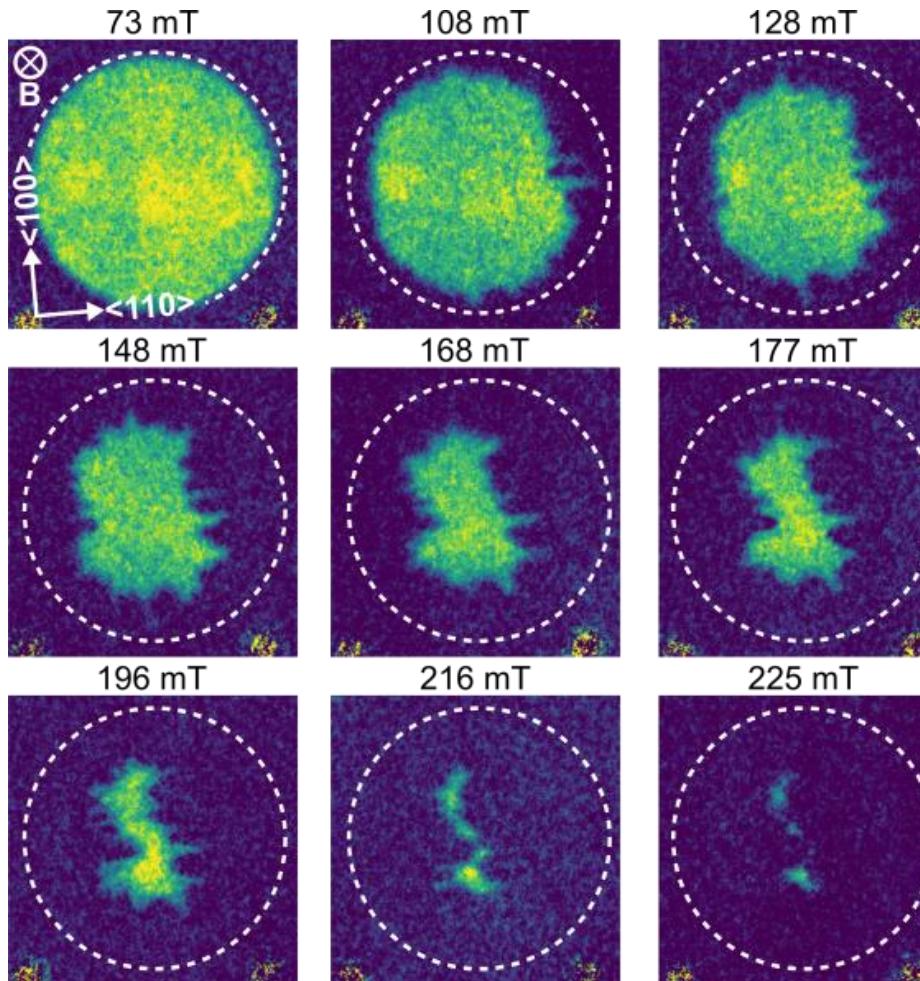
Field Cooling into the IMS



What changes?

- domain size
- domain density
- scattering contrast

The Intermediate Mixed State in Niobium



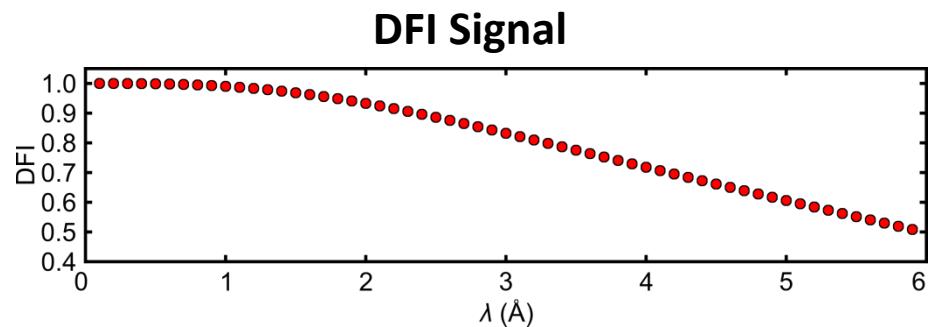
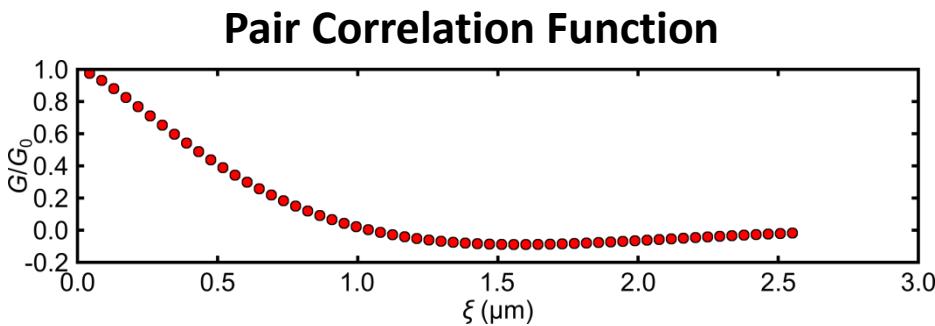
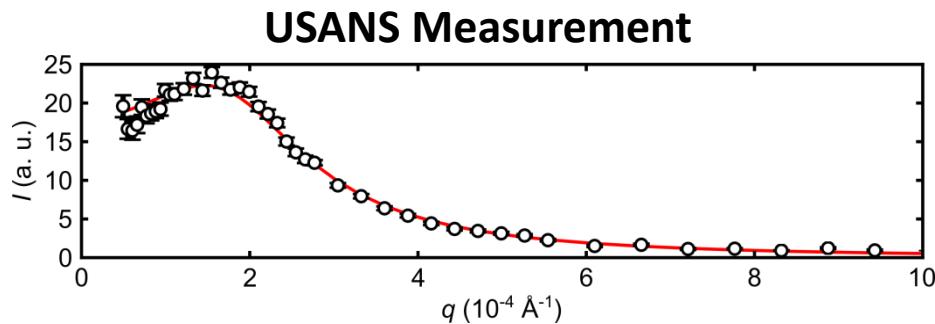
Additional field enters the sample

- IMS in the center
- FLL at the edge
- pinning along crystal axes

Heterogeneous Phase transition

T. Reimann, Ph.D Thesis (2016)

From USANS to nGI



$$\left(\frac{d\sigma(q_x)}{d\Omega} \right)_{\text{slit}}$$



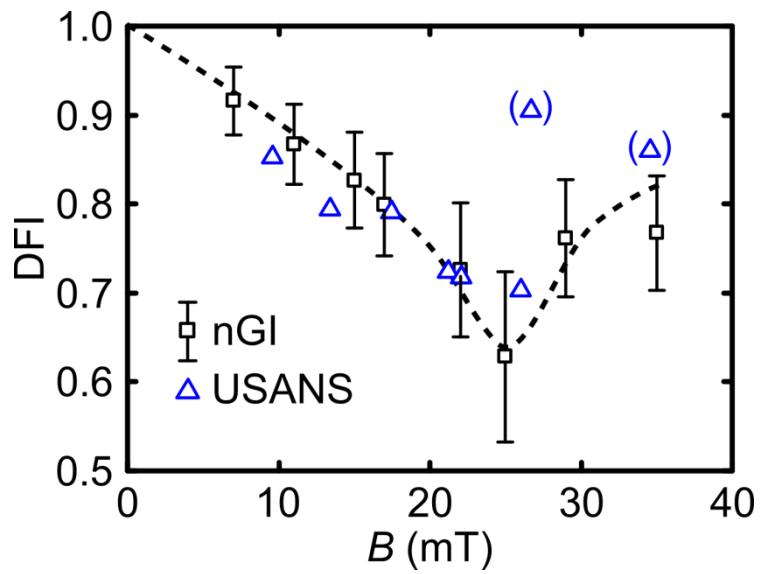
$$G(x, 0)$$

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T. Reimann, Ph.D Thesis (2016)

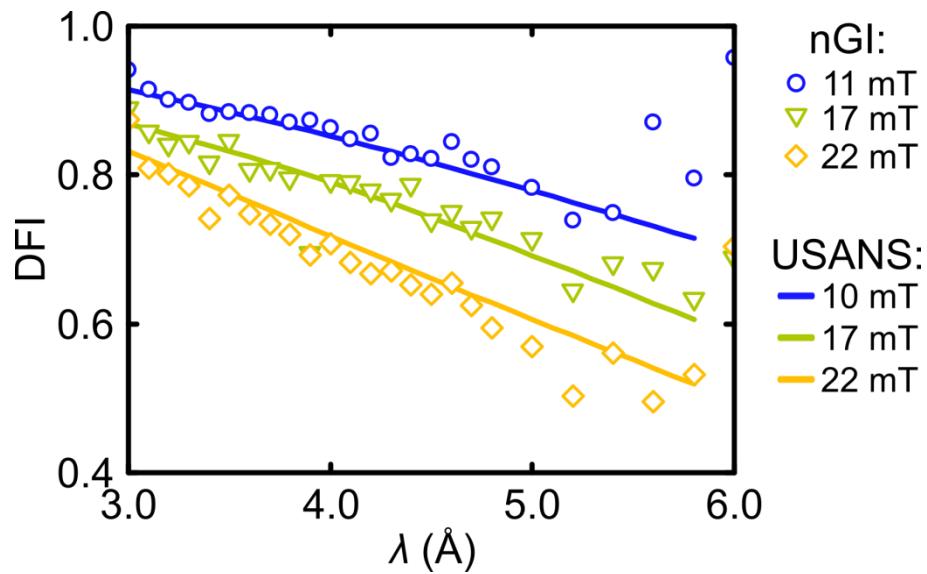
Magnetic Field Scan

Domain size changes
(USANS result)



Wavelength Scan

Probed ξ_{GI} changes



T. Reimann, Ph.D Thesis (2016)